

Séminaire - FIME

Probabilistic representation
of a class of nonconservative nonlinear PDE

ANTHONY LECAVIL¹, NADIA OUDJANE² AND
FRANCESCO RUSSO³.

June, 5th 2015

-
1. ENSTA-ParisTech, UMA
 2. EDF R&D, FiME
 3. ENSTA-ParisTech, UMA

Summary

- 1 General framework
 - Motivations
 - State of the art
 - Statement of the problem
 - Main results
- 2 Existence and uniqueness of the new Nonlinear SDEs
 - Existence and uniqueness under weaker assumptions
 - Link with the Partial Integro-Differential Equation
- 3 Numerical approximation scheme
 - Particle system and Propagation of chaos
 - Time discretization scheme
 - Simulations results

Plan

- 1 General framework
 - Motivations
 - State of the art
 - Statement of the problem
 - Main results

We consider the following **non conservative** and **nonlinear** PDE

$$\left\{ \begin{array}{l} \partial_t u = \frac{1}{2} \sum_{i,j=1}^d \partial_{ij}^2 ((\Phi \Phi^t)_{i,j}(t, x, u) u) - \operatorname{div} (g(t, x, u) u) + \Lambda(t, x, u) u \\ u(0, dx) = \zeta_0(dx) . \end{array} \right.$$

- Aim 1 : Find a **forward probabilistic representation** of the PDE
- Aim 2 : Propose a **numerical approximation** of the solution which is both
 - ① *less sensitive to the dimension as a Monte Carlo scheme ;*
 - ② *able to concentrate the computing efforts in the region of interest as a forward representation.*

Major contributions since the sixties

- **Conservative PDE** : $\int_{\mathbb{R}^d} u_t(x) dx = 1$ for all $t \in [0, T]$

$$\partial_t u_t = \frac{1}{2} \partial_{xx}^2 (\Phi(x, u_t) u_t) - \partial_x (b(x, u_t) u_t), \quad (\Lambda = 0) \quad \text{where}$$

$$\begin{cases} \Phi(x, u_t) & := \int_{\mathbb{R}^d} K^\Phi(x, y) u_t(dy), \\ g(x, u_t) & := \int_{\mathbb{R}^d} K^g(x, y) u_t(dy), \end{cases}$$

Integral dependence on u and not point dependence on u .

- McKean introduced the notion of **nonlinear SDE (NLSDE)**

$$\begin{cases} Y_t = Y_0 + \int_0^t \Phi(Y_s, u_s) dW_s + \int_0^t g(Y_s, u_s) ds \\ u_t \text{ is the density of the law of } Y_t, \end{cases} \quad (1.1)$$

- Propose an **interacting particle system (IPS)** whose the limit is a sol. of PDE : **propagation of chaos** estimates.

- Méléard et al. have studied, under **smooth assumptions**, exist./uniqu. of

$$\begin{cases} Y_t = Y_0 + \int_0^t \Phi(u(s, Y_s)) dW_s + \int_0^t g(u(s, Y_s)) ds \\ u_t \text{ is the density of the law of } Y_t \end{cases} \quad (1.2)$$

\implies **point dependence** on u , i.e. $K^\Phi(\cdot, y) = K^g(\cdot, y) = \delta_y$.

- They also proved that the **regularized version**

$$\begin{cases} Y_t^\varepsilon = Y_0 + \int_0^t \Phi((K_\varepsilon * u^\varepsilon)(s, Y_s^\varepsilon)) dW_s + \int_0^t g((K_\varepsilon * u^\varepsilon)(s, Y_s^\varepsilon)) ds \\ u_t^\varepsilon \text{ is the density of the law of } Y_t^\varepsilon \end{cases}$$

strongly converges to (1.2) when $K_\varepsilon \xrightarrow{\varepsilon \rightarrow 0} \delta$.

- Benachour et al. have proved exist./uniq. of

$$\begin{cases} Y_t = Y_0 + \int_0^t \Phi(u(s, Y_s)) dW_s \\ u_t \text{ is the density of the law of } Y_t, \end{cases} \quad (1.3)$$

with $\Phi : x \in \mathbb{R} \mapsto x^{\frac{k-1}{2}}$, $k \geq 1$.

Russo et al. have extended (1.3) for Φ only **bounded and measurable**.

- This representation is associated to the Porous Media Equation

$$\partial_t u = \frac{1}{2} \partial_{xx}^2 (u \Phi^2(u)).$$

- Framework : **Nonconservative nonlinear PDE** of the form

$$\begin{cases} \partial_t u = \frac{1}{2} \sum_{i,j=1}^d \partial_{ij}^2 ((\Phi \Phi^t)_{i,j}(t, x, u) u) - \operatorname{div} (g(t, x, u) u) + \Lambda(t, x, u) u \\ u(0, dx) = \zeta_0(dx), \end{cases}$$

where

- ζ_0 is a probability measure on \mathbb{R}^d ;
- $\Phi : [0, T] \times \mathbb{R}^d \times \mathbb{R} \rightarrow \mathbb{R}^{d \times d}$, $g : [0, T] \times \mathbb{R}^d \times \mathbb{R} \rightarrow \mathbb{R}^d$,
 $\Lambda : [0, T] \times \mathbb{R}^d \times \mathbb{R} \rightarrow \mathbb{R}$ are bounded and measurable functions ;
- $u(0, dx) = \zeta_0(dx)$ means $u(t, x) dx \xrightarrow[t \rightarrow 0]{} \zeta_0(dx)$ weakly

Nonconservative $\iff \int_{\mathbb{R}^d} u(t, x) dx = fct(t) \iff \Lambda \neq 0$.

- Our idea : consider the following representation

$$\left\{ \begin{array}{l} Y_t = Y_0 + \int_0^t \Phi(s, Y_s, \mathbf{u}(s, Y_s)) dW_s + \int_0^t g(s, Y_s, \mathbf{u}(s, Y_s)) ds \\ \mathbf{u}(t, \cdot) := \frac{d\nu_t}{dx} \quad \text{such that for any } \varphi \in \mathcal{C}_b(\mathbb{R}^d, \mathbb{R}) \\ \nu_t(\varphi) := \mathbb{E} \left[\varphi(Y_t) \exp \left\{ \int_0^t \Lambda(s, Y_s, \mathbf{u}(s, Y_s)) ds \right\} \right], \end{array} \right.$$

Observations :

- $\int_{\mathbb{R}^d} u(t, x) dx = \mathbb{E} \left[\exp \left\{ \int_0^t \Lambda(s, Y_s, \mathbf{u}(s, Y_s)) ds \right\} \right]$.
- The measure ν_t needs the **law of all the process** Y ($\in \mathcal{P}(\mathcal{C}([0, T], \mathbb{R}^d))$) and not only marginals laws.
- **point dependence** on u in Φ and $g \Rightarrow$ technical difficulty.

- Bypass the difficulty : consider a *regularized version of NLSDE*,

$$\begin{cases} Y_t = Y_0 + \int_0^t \Phi(s, Y_s, \mathbf{u}(s, Y_s)) dW_s + \int_0^t g(s, Y_s, \mathbf{u}(s, Y_s)) ds \\ \mathbf{u}(t, y) = \mathbb{E} \left[K(y - Y_t) \exp \left\{ \int_0^t \Lambda(s, Y_s, \mathbf{u}(s, Y_s)) ds \right\} \right]. \end{cases}$$

- **Integral dependence** on $\mathcal{L}(Y_\cdot) \in \mathcal{P}(\mathcal{C}^d)$.
- **u depends on itself** \implies main difference with the cases already covered in the literature.
- Formally, $\Lambda = 0$ and $K = \delta$: cases already developed by Méléard and al. (i.e. conservative case).

Main results of existence and uniqueness

1 "Lipschitz" case : If

- ζ_0 admits a 2nd order moment,
- Φ, g, Λ are bounded, **uniformly Lipschitz w.r.t. t** ,

there is a unique **strong solution** (Y, u) .

2 "Semi-weak" case : If

- ζ_0 admits a 2nd order moment,
- Φ, g are bounded and **uniformly Lipschitz w.r.t. t** ,
- Λ is only **continuous**,

there is a **(non-unique) strong solution** (Y, u) .

3 "Weak" case : If

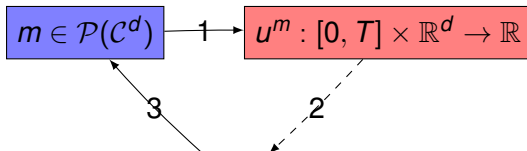
- Φ, g, Λ are bounded and **continuous**

there is a **weak solution** (Y, u) .

Plan

- 1 General framework
 - Motivations
 - State of the art
 - Statement of the problem
 - Main results
- 2 Existence and uniqueness of the new Nonlinear SDEs
 - Existence and uniqueness under weaker assumptions
 - Link with the Partial Integro-Differential Equation
- 3 Numerical approximation scheme
 - Particle system and Propagation of chaos
 - Time discretization scheme
 - Simulations results

- Existence/uniqueness of (Y, u) in *Lipschitz case* relies on :



$$Y_t^m = Y_0 + \int_0^t \Phi(s, Y_s, u^m(s, Y_s)) dW_s + \int_0^t g(s, Y_s, u^m(s, Y_s)) ds$$

where we recall that

$$u^m(t, y) = \mathbb{E} \left[K(y - Y_t) e^{\int_0^t \Lambda(s, Y_s, u^m(s, Y_s)) ds} \right], \text{ with } m := \mathcal{L}(Y).$$

- Aim 1 : Existence / uniqueness of the map $m \mapsto u^m$.
- Aim 2 : Existence / uniqueness of the map $\Theta : m \mapsto \mathcal{L}(Y^m)$.

- Aim 1 : Existence and uniqueness of u^m for fixed $m \in \mathcal{P}(C^d)$

Does it exist a unique function $u : [0, T] \times \mathbb{R}^d \rightarrow \mathbb{R}$ such that for
all $(t, y) \in [0, T] \times \mathbb{R}^d$,

$$u(t, y) = \int_{C^d} K(y - X_t(\omega)) \exp \left\{ \int_0^t \Lambda(s, X_s(\omega), u(s, X_s(\omega))) ds \right\} dm(\omega)$$

Yes, if Λ is bounded and uniformly Lipschitz w.r.t. t .

- Idea of the proof : fixed-point argument.

Remark

- *Existence and uniqueness of u is obtained for all $m \in \mathcal{P}(\mathcal{C}^d)$.*
- *Only the hypothesis on Λ are used here(= bounded and uniformly Lipschitz w.r.t. t) and not those of Φ, g .*
- *Uniqueness is lost if Λ is only continuous !!!*

Stability properties for $u^m(t, y) := u(m, t, y)$ under various norms :

• $\forall (m, m') \in \mathcal{P}(\mathcal{C}^d) \times \mathcal{P}(\mathcal{C}^d), \forall (t, y, y') \in [0, T] \times \mathbb{R}^d \times \mathbb{R}^d :$

$$|u^m(t, y) - u^{m'}(t, y')|^2 \leq \mathfrak{C}_{K, \Lambda}(T) \left[|y - y'|^2 + |\widetilde{W}_t(m, m')|^2 \right],$$

where the map

$$(m, m') \in \mathcal{P}(\mathcal{C}^d) \times \mathcal{P}(\mathcal{C}^d) \mapsto |\widetilde{W}_T(m, m')|^2$$

is the 2-Wasserstein distance on the space of Borel probability measures on \mathcal{C}^d , s.th. for all $t \in [0, T]$,

$$|\widetilde{W}_t(m, m')|^2 := \inf_{\mu \in \widetilde{\Pi}(m, m')} \int_{\mathcal{C}^d \times \mathcal{C}^d} \left(\sup_{0 \leq s \leq t} |X_s(\omega) - X_s(\omega')|^2 \wedge 1 \right) d\mu(\omega, \omega')$$

- The function

$$(m, t, x) \mapsto u^m(t, x)$$

is continuous on $\mathcal{P}(\mathcal{C}^d) \times [0, T] \times \mathbb{R}^d$ where $\mathcal{P}(\mathcal{C}^d)$ is endowed with the topology of weak convergence.

- Suppose here that $K \in W^{1,2}(\mathbb{R}^d)$.

For any $t \in [0, T]$, $(m, m') \in \mathcal{P}_2(\mathcal{C}^d) \times \mathcal{P}_2(\mathcal{C}^d)$,

$$\|u^m(t, \cdot) - u^{m'}(t, \cdot)\|_2^2 \leq \tilde{\mathfrak{C}}_{K,\Lambda}(T) |W_t(m, m')|^2,$$

where $\|\cdot\|_2$ is the standard $L^2(\mathbb{R}^d)$ or $L^2(\mathbb{R}^d, \mathbb{R}^d)$ -norms.

- Suppose (additionally) that $\mathcal{F}(K) \in L^1(\mathbb{R}^d)$. Then
 $\exists \bar{c}_{K,\Lambda}(t) > 0$ for all $(m, t) \in \mathcal{P}(C^d) \times [0, T]$,

$$\mathbb{E}[\|u^{S^N(\xi)}(t, \cdot) - u^m(t, \cdot)\|_\infty^2] \leq \bar{c}_{K,\Lambda}(T) \sup_{\substack{\varphi \in C_b(C^d) \\ \|\varphi\|_\infty \leq 1}} \mathbb{E}[|\langle S^N(\xi) - m, \varphi \rangle|^2]$$

where

$$S^N(\xi) := \frac{1}{N} \sum_{i=1}^N \delta_{\xi^i}$$

for $(\xi^i, 1 \leq i \leq N)$ given continuous processes.

- Aim 2 : Existence and uniqueness of the process Y
 - The map $u : m \mapsto u^m$ is defined for all $m \in \mathcal{P}(\mathcal{C}^d)$;
 - We apply a fixed-point argument (see Sznitman) to the map

$$\Theta : \mathcal{P}_2(\mathcal{C}^d) \rightarrow \mathcal{P}_2(\mathcal{C}^d) ,$$

defined by $\Theta(m) = \mathcal{L}(Y^m)$ s.th.

$$Y_t^m = Y^0 + \int_0^t \Phi(s, Y_s^m, u^m(s, Y_s^m)) dW_s + \int_0^t g(s, Y_s^m, u^m(s, Y_s^m)) ds$$

with $\mathcal{P}_2(\mathcal{C}^d)$ equipped with 2-Wasserstein distance W_T^2 .

Plan

- 1 General framework
 - Motivations
 - State of the art
 - Statement of the problem
 - Main results
- 2 Existence and uniqueness of the new Nonlinear SDEs
 - Existence and uniqueness under weaker assumptions
 - Link with the Partial Integro-Differential Equation
- 3 Numerical approximation scheme
 - Particle system and Propagation of chaos
 - Time discretization scheme
 - Simulations results

- Existence in *Semi-weak case* and *Weak case* :

$$\begin{cases} Y_t = Y_0 + \int_0^t \Phi(s, Y_s, u(s, Y_s)) dW_s + \int_0^t g(s, Y_s, u(s, Y_s)) ds \\ u(t, y) = \mathbb{E} \left[K(y - Y_t) \exp \left\{ \int_0^t \Lambda(s, Y_s, u(s, Y_s)) ds \right\} \right], \end{cases}$$

admits a solution in semi-weak and weak case.

The proof consists in

- 1 **regularizing** the coefficients Φ , g , Λ with a mollifier $(\varphi_n)_{n \in \mathbb{N}}$.
- 2 using the **Lipschitz / Semi-weak case result** for mollified coefficients \implies existence of $(Y^n, u^n)_{n \in \mathbb{N}}$.
- 3 **convergence of $(u^n)_n$** and identification of the limit.
- 4 identify the limit of (Y^n) (stability of SDEs / martingale formulation).

Plan

- 1 **General framework**
 - Motivations
 - State of the art
 - Statement of the problem
 - Main results
- 2 **Existence and uniqueness of the new Nonlinear SDEs**
 - Existence and uniqueness under weaker assumptions
 - **Link with the Partial Integro-Differential Equation**
- 3 **Numerical approximation scheme**
 - Particle system and Propagation of chaos
 - Time discretization scheme
 - Simulations results

- Link with PIDE : **Ito's formula** implies that (Y, u) solution of (regularized) NLSDE is related to the partial integro-differential equation (PIDE)

$$\left\{ \begin{array}{l} \partial_t v = \frac{1}{2} \sum_{i,j=1}^d \partial_{ij}^2 ((\Phi \Phi^t)_{i,j}(t, x, K * v) v) - \operatorname{div} (g(t, x, K * v) v) \\ + \Lambda(t, x, K * v) v \\ v_0 = \zeta_0, \end{array} \right.$$

by the relation

$$v_t(\cdot) = (K * v_t)(\cdot) = \int_{R^d} K(\cdot - y) v_t(dy) .$$

Plan

- 1 General framework
 - Motivations
 - State of the art
 - Statement of the problem
 - Main results
- 2 Existence and uniqueness of the new Nonlinear SDEs
 - Existence and uniqueness under weaker assumptions
 - Link with the Partial Integro-Differential Equation
- 3 Numerical approximation scheme
 - Particle system and Propagation of chaos
 - Time discretization scheme
 - Simulations results

- In all the sequel, only assumptions of *Lipschitz case* will be satisfied.

- Interacting Particle System (IPS)

For fixed i.i.d. r.v. $(Y_0^i)_{i=1, \dots, N}$ and $(W^i)_{i=1, \dots, N}$ a family of independent Brownian motions, the IPS $\xi := (\xi^{i,N})_{i=1, \dots, N}$ is defined by

$$\begin{cases} \xi_t^{i,N} = Y_0^i + \int_0^t \Phi_s(\xi_s^{i,N}, u_s^{S^N(\xi)}(\xi_s^{i,N})) dW_s^i + \int_0^t g_s(\xi_s^{i,N}, u_s^{S^N(\xi)}(\xi_s^{i,N})) ds \\ u_t^{S^N(\xi)}(x) = \frac{1}{N} \sum_{j=1}^N K(x - \xi_t^{j,N}) \exp\left(\int_0^t \Lambda(r, \xi_r^{j,N}, u_r^{S^N(\xi)}(\xi_r^{j,N})) dr\right), \end{cases}$$

with $S^N(\xi) := \frac{1}{N} \sum_{j=1}^N \delta_{\xi^{j,N}}$, empirical measure associated to ξ .

For such systems, **propagation of chaos** \equiv "asymptotic independence" of the components $(\xi^i)_{i=1, \dots, N}$ when the size N (=number of components) goes to $+\infty$.

Main ideas :

- Transform a d -dimensional (regularized) NLSDE into a $d \times N$ -dimensional classical SDEs.
- The function $u^{S^N(\xi)}$ can be seen as the "mixing/interaction term". It can be written

$$u_t^{S^N(\xi)}(x) = F(t, x, \xi_t^1, \dots, \xi_t^N, \underbrace{(\xi_{\cdot \wedge t}^1), \dots, (\xi_{\cdot \wedge t}^N)}_{\text{past of the trajectories}}).$$

- Dimension of the state space ($=(\mathbb{R}^d)^N$) depends on $N \neq \omega \mapsto S^N(\xi(\omega)) \in \mathcal{P}(\mathcal{C}^d)$ with $\dim(\mathcal{P}(\mathcal{C}^d)) = \infty$.

- If $\xi := (\xi^i)_{i=1, \dots, N}$ are i.i.d. \mathbb{R}^d -valued r.v. according to $\mu \in \mathcal{P}(\mathbb{R}^d)$,

$$\langle S^N(\xi), \varphi \rangle = \frac{1}{N} \sum_{i=1}^N \varphi(\xi^i) \xrightarrow[N \rightarrow +\infty]{p.s.} \langle \mu, \varphi \rangle,$$

by the Strong law of large numbers.

Existence/uniqueness of such IPS

- Non-anticipative property : $\forall (s, x) \in [0, T] \times \mathbb{R}^d$,

$$u_s^{S^N(\xi)}(x) = u_s^{S^N((\xi_r, 0 \leq r \leq s))}(x).$$

\implies all integrands of IPS are adapted and so, Itô's integral is well-defined.

- Lipschitz property of integrands :

Lipschitz property of $m \mapsto u^m$ implies that

$$(s, \bar{\xi}) \in [0, T] \times (\mathcal{C}^d)^N \mapsto \Phi(s, \bar{\xi}_s^{i,N}, u_s^{S^N(\bar{\xi})}(\bar{\xi}_s^{i,N}))$$

and

$$(s, \bar{\xi}) \in [0, T] \times (\mathcal{C}^d)^N \mapsto g(s, \bar{\xi}_s^{i,N}, u_s^{S^N(\bar{\xi})}(\bar{\xi}_s^{i,N}))$$

are Lipschitz.

Consequently, classical results for **path-dependent SDEs** give existence/uniqueness.

Coupling technique :

Let $(Y^i)_{i=1, \dots, N}$ be solutions of

$$\begin{cases} Y_t^i = Y_0^i + \int_0^t \Phi(s, Y_s^i, u^{m_i}(s, Y_s^i)) dW_s^i + \int_0^t g(s, Y_s^i, u^{m_i}(s, Y_s^i)) ds \\ u^{m_i}(t, x) = \mathbb{E} \left[K(x - Y_t^i) \exp \left(\int_0^t \Lambda(r, Y_r^i, u^{m_i}(s, Y_s^i)) \right) \right], \end{cases}$$

where $(W^i)_{i=1, \dots, N}$ is the same family of independent Brownian motions driving the IPS $(\xi^{i,N})_{i=1, \dots, N}$. Then,

- (Y^1, \dots, Y^N) are i.i.d. and their common law will be denoted by m^0 .

Theorem

Under some assumptions, the following inequalities hold :

$$\begin{aligned}\mathbb{E}[\|u_t^{S^N(\xi)} - u_t^{m_0}\|_\infty^2] &\leq \frac{C}{N} \\ \mathbb{E}[\sup_{0 \leq s \leq t} |\xi_s^{i,N} - Y_s^i|^2] &\leq \frac{C}{N} \\ \mathbb{E}[\|u_t^{S^N(\xi)} - u_t^{m_0}\|_2^2] &\leq \frac{C}{N},\end{aligned}$$

for all $i \in \{1, \dots, N\}$ and where C does not depend on N .

Time discretization version of the IPS

To simplify notations, we set $g \equiv 0$.

- Euler Scheme : for $k = 1, \dots, n$

$$\left\{ \begin{array}{l} \tilde{\xi}_{t_{k+1}}^{i,N} = \tilde{\xi}_{t_k}^{i,N} + \Phi(t_k, \tilde{\xi}_{t_k}^{i,N}, \tilde{\mathbf{v}}_{t_k}(\tilde{\xi}_{t_k}^{i,N}))\mathcal{N}(0, \delta t) \\ \tilde{\xi}_0^{i,N} = Y_0^i \\ \tilde{\mathbf{v}}_{t_{k+1}}(y) = \frac{1}{N} \sum_{j=1}^N K(y - \tilde{\xi}_{t_{k+1}}^{j,N}) e^{\left\{ \sum_{p=0}^k \Lambda(t_p, \tilde{\xi}_{t_p}^{j,N}, \tilde{\mathbf{v}}_{t_p}(\tilde{\xi}_{t_p}^{j,N})) \delta t \right\}} \end{array} \right. ,$$

where $0 \leq t_0 < \dots < t_k = k * \delta t < \dots < t_n \leq T$ is a regular time grid.

Theorem

Under Lipschitz continuity assumption, we have

$$\mathbb{E}[\|\tilde{v}_t - u_t^{S^N(\xi)}\|_\infty^2] + \sup_{i=1, \dots, N} \mathbb{E} \left[\sup_{s \leq t} |\tilde{\xi}_s^{i,N} - \xi_s^{i,N}|^2 \right] \leq C_{K, \Lambda, T} \delta t .$$

and the Mean Integrated Squared Error (MISE) verifies

$$\mathbb{E}[\|\tilde{v}_t - u_t^{S^N(\xi)}\|_2^2] \leq C_{K, \Lambda, T} \delta t .$$

Remark

- The constant $C_{K,\Lambda,T}$ does not depend on N and δt .
- By the two previous Theorems, we have for all $t \in [0, T]$

$$\mathbb{E} \left[\|\tilde{\mathbf{v}}_t - u_t^{m^0}\|_\infty^2 \right] \leq C \left(\delta t + \frac{1}{N} \right).$$

Initialization for $k = 0$

- ① Generate $(\xi_0^i)_{i=1, \dots, N}$ i.i.d. $\sim v(0, x) dx$;
- ② Set $G_0^i = 1, i = 1, \dots, N$;
- ③ Set $\tilde{v}_0(\cdot) := v(0, \cdot)$;

Iterations for $k = 1, \dots, n-1$

- Independently for each particle $\tilde{\zeta}_k^{j,N}$ for $j = 1, \dots, N$,

$$\tilde{\zeta}_{k+1}^{j,N} = \tilde{\zeta}_k^{j,N} + \Phi(t_k, \tilde{\zeta}_k^{j,N}, \tilde{v}_k(\tilde{\zeta}_k^{j,N})) \mathcal{N}(0, \delta t)$$

- Set

$$G_{k+1}^j := G_k^j \times \exp \left(\Lambda(t_k, \tilde{\zeta}_k^{j,N}, \tilde{v}_k(\tilde{\zeta}_k^{j,N})) \delta t \right);$$

- Set $\tilde{v}_k(\cdot) = \frac{1}{N} \sum_{i=1}^N G_k^i \times K_h(\cdot - \tilde{\zeta}_{k-1}^{i,N})$

- ***In all the sequel, we expose an empirical analysis for which the assumptions of the theorems are not necessarily satisfied.***
- ***Since***

$$\mathbb{E} \left[\|\tilde{v}_t - u_t^{S^N(\xi)}\|_\infty^2 \right] \leq C \delta t$$

where $C := C(\|K\|_\infty, \|\Lambda\|_\infty, L_K, L_\Lambda, \|\nabla K\|_2, T)$, notice that we neglect the time discretization in the present empirical analysis.

Aim : show how the particle system can be used to estimate u ,
 solution of the PDE

$$\left\{ \begin{array}{l} \partial_t u = \frac{1}{2} \sum_{i,j=1}^d \partial_{ij}^2 ((\Phi \Phi^t)_{i,j}(t, x, u)u) - \operatorname{div}(g(t, x, u)u) + \Lambda(t, x, u)u \\ u(0, dx) = \zeta_0(dx) . \end{array} \right.$$

Let us consider the interacting particle system $\xi^{i,N,\varepsilon}$, where
 $K = K_\varepsilon$ for $\varepsilon > 0$.

$$\begin{aligned} \xi_t^{i,N,\varepsilon} &= Y_0^i + \int_0^t \Phi_s(\xi_s^{i,N,\varepsilon}, u_s^{SN(\xi^\varepsilon)}(\xi_s^{i,N,\varepsilon})) dW_s^i \\ &\quad + \int_0^t g_s(\xi_s^{i,N,\varepsilon}, u_s^{SN(\xi^\varepsilon)}(\xi_s^{i,N,\varepsilon})) ds , \end{aligned}$$

where

$$u_t^{S^N(\xi^\varepsilon)}(x) = \frac{1}{N} \sum_{j=1}^N K_\varepsilon(x - \xi_t^{j,N,\varepsilon}) e^{\int_0^t \Lambda(r, \xi_r^{j,N,\varepsilon}, u_r^{S^N(\xi^\varepsilon)}(\xi_r^{j,N,\varepsilon})) dr}.$$

We are going to try to show this empirically. To this end, we consider the Mean Integrated Squared Error (MISE) that we decompose as the **Variance** and the **Bias**,

$$\begin{aligned} MISE_t(\varepsilon, N) &:= V_t(\varepsilon, N) + B^2(\varepsilon, N) \\ &= \mathbb{E} \left[\|u_t^{S^N(\xi^\varepsilon)} - \mathbb{E}[u_t^{S^N(\xi^\varepsilon)}]\|_2^2 \right] + \mathbb{E} \left[\|\mathbb{E}[u_t^{S^N(\xi^\varepsilon)}] - v_t\|_2^2 \right]. \end{aligned}$$

Ideally, we would like that

$$\mathbb{E}[u_t^{S^N(\xi^\varepsilon)}] \sim u_t^{m^0, \varepsilon}.$$

If the **propagation of chaos** holds, it means that the particles $\xi^{N,\varepsilon}$ are **close to an i.i.d. system** according to m^0 , which is the common law of the processes Y^i , $1 \leq i \leq N$. Then, in the particular case where $\Lambda(t, x, u) := \Lambda(t, x)$,

$$\begin{aligned} \mathbb{E}[u_t^{S^N(\xi^\varepsilon)}] &= \frac{1}{N} \mathbb{E} \left[\sum_{j=1}^N K(\cdot - \xi_t^{j,N,\varepsilon}) \exp \left(\int_0^t \Lambda(r, \xi^{j,N,\varepsilon}) dr \right) \right] \\ &\approx \mathbb{E} \left[K_\varepsilon(\cdot - Y_t^{1,\varepsilon}) \exp \left(\int_0^t \Lambda(r, Y^{1,\varepsilon}) dr \right) \right] \\ &= u_t^{m^0, \varepsilon} \end{aligned}$$

We expect to observe the same behavior for the case $\Lambda = \Lambda(t, x, u)$ depends on u .

Finally,

$$V_t(\varepsilon, N) \approx \mathbb{E} \left[\|u_t^{S^N(\xi^\varepsilon)} - u_t^{m^0, \varepsilon}\|_2^2 \right]$$

and

$$B_t^2(\varepsilon, N) \approx B_t^2(\varepsilon) := \mathbb{E} \left[\|u_t^{m^0, \varepsilon} - v_t\|_2^2 \right] .$$

In other words,

$V_t(\varepsilon, N)$ corresponds to the convergence of the particles system (i.e. when $N \rightarrow +\infty$, for fixed $\varepsilon > 0$) ,

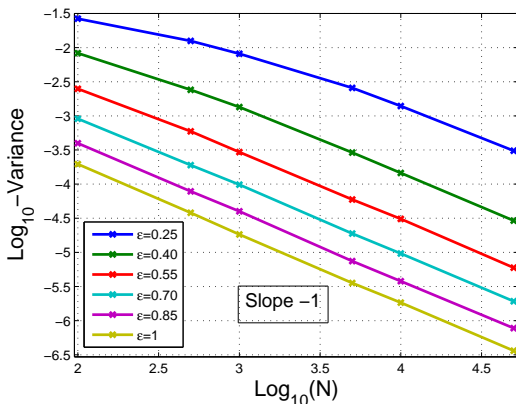
and

$B_t^2(\varepsilon, N)$ corresponds to the convergence of the regularized NLSDE (i.e. when $\varepsilon \rightarrow 0$) .

Variance Analysis (1) : Behavior w.r.t. N

Simulations given below for $d = 5$, $T = 1$ give us

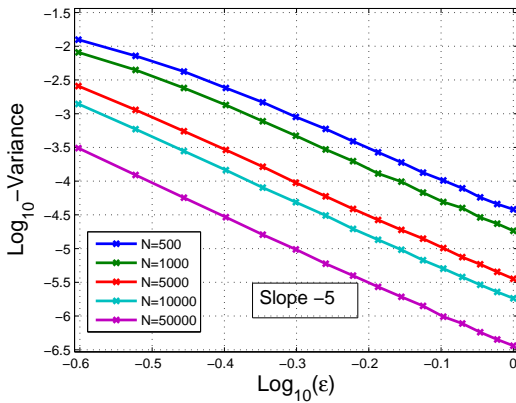
$$V_t(\varepsilon, N) \sim \frac{1}{N\varepsilon^d}$$



Variance Analysis (2) : Behavior w.r.t. ε :

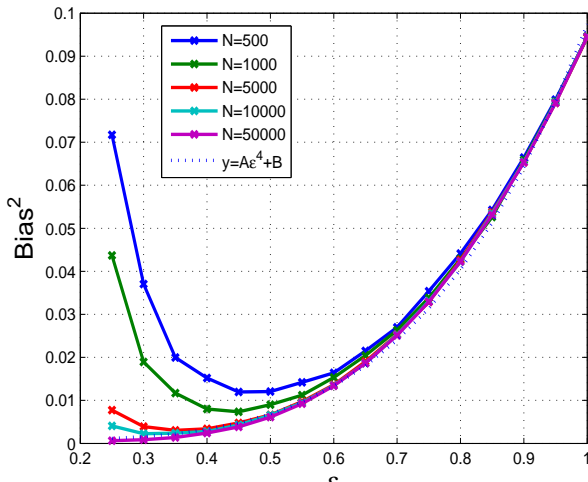
Simulations given below for $d = 5$, $T = 1$ give us

$$V_t(\varepsilon, N) \sim \frac{1}{N\varepsilon^d}$$



Bias Analysis (2) :

Simulations give us $B^2(\varepsilon) \sim \varepsilon^4$



Thank you for your attention

References I